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# **CPA approach: On Heavy Carbon Doping of MgB2**



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# Outline

- Experimental data
- Calculational methods
- Smeared band structure
- CPA & Supercell calculation results
- VCA results
- $\diamond$   $\sigma$  hole concentration

## Onclusion

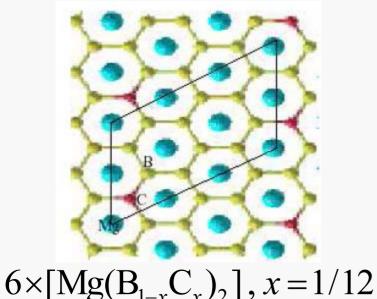
# **Experimental data**

- MgB<sub>2</sub>: a real possibility for a high field conductor, due to high parallel and perpendicular critical fields
- $\otimes Mg_{1-v}Al_vB_2$ :
  - 1. up to y=0.10, T<sub>c</sub> decreases smoothly
  - 2. 0.10 < y < 0.25, two phase behavior (two *c* parameters)
  - 3. y=0.25, back to single phase with vanishing T<sub>c</sub>
- $\otimes Mg(B_{1-x}C_x)_2$  :
  - 1.  $dT_c / dx = 1K / \% C$
  - 2. H<sub>c2</sub> increases strongly with C content. (~33-35 T)
  - 3. It remains the 2-band gap superconductor for  $10\pm2\%$  C substitution for B.

# **Calculational Method**

#### LAPW (Wien2K)

- P.Blaha et al., Comput. Phys. Commun. <u>59</u>, 399 (1990)
- VCA (10% C doping)
- 2x3 Supercell method: ordered impurity calculation



### FPLO

K. Koepernik and H. Eschrig, PRB <u>59</u>, 1743 (1999) http://www.ifw-dresden.de/agtheo/FPLO/

 Coherent Potential Approximation (CPA): disordered (randomly substituted) alloy calculations

K. Koepernik, B. Velicky, R. Hayn, and H. Eschrig, PRB <u>55</u>, 5717 (1997)

# **CPA detail**

### Solution & Goal

- 1) determine the filling of the  $\sigma$  band hole states
- 2) obtain the broadening (and potentially splitting) of bands
- ♦ Doping level:  $Mg(B_{1-x}C_x)_2$

x	Goal
0.0001	"perfect crystal" reference for evaluation of the numerical algorithms in CPA
0.0833	to compare with $x=1/12$ supercell calculation
0.10 & 0.20	Representative of the system toward the achievable upper range

- *k*-mesh: 45x45x10 (1152 irreducible *k*-point)
- Valence orbitals: 2s2p3s3p3d (Mg), 2s2p3d (B&C)

## Smeared band structure (Bloch spectral density)

For ordered single particle spectra,

$$A_{k}^{bl}(\omega) = \sum_{n} \delta(\omega - \varepsilon_{kn})$$
  
(*if*  $\omega = \varepsilon_{kn}$ , black dot  
*otherwise*, white)  $\Rightarrow$  usual band struture

#### In disorder case,

$$A_{k}^{bl}\left(\omega\right) = Tr\left(S\left(k\right)\Gamma\left(k,\omega\right)\right)$$

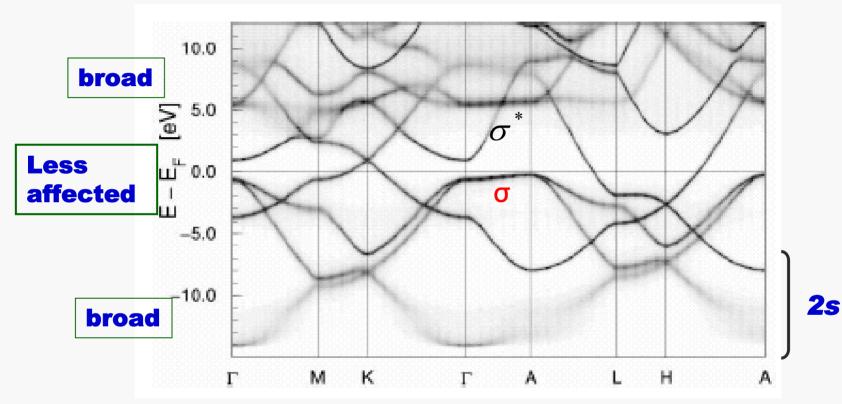
where S(k): overlap matrix

 $\Gamma(k,\omega)$ : nonstochastic nonlocal coherent Green function

Iarge A : dense darkness small A: light

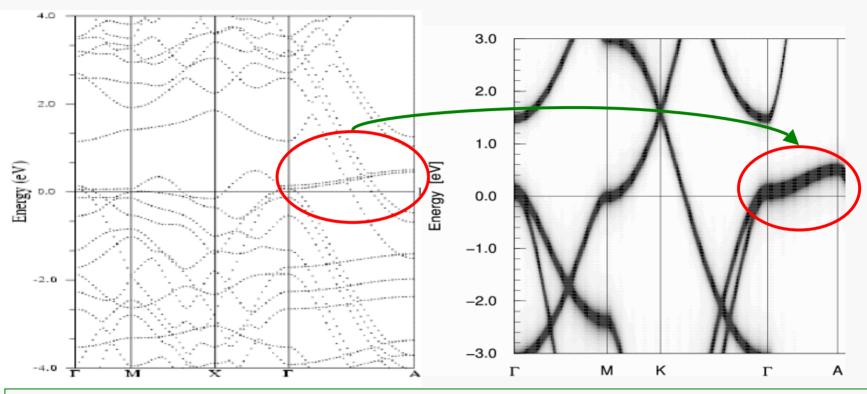
intermediate region: continuous gray gradients

## **CPA** | (x = 0.20)



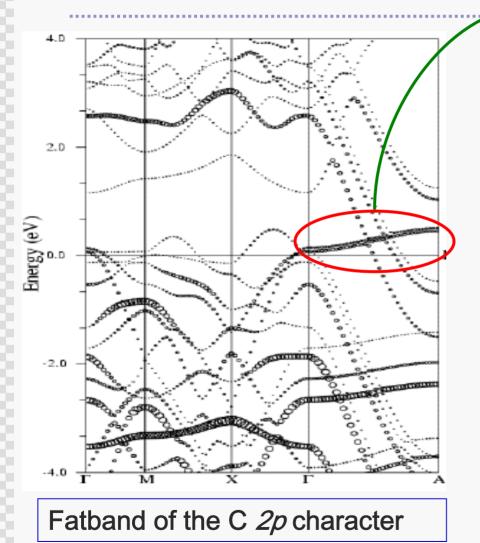
- Largest broadening (2s): C 2s state is noticeably lower in energy leading to increased smearing.
- ♦ Near E<sub>F</sub>, less affected by the chemical disorder.
- ♦ The σ band holes are completely filled.
   ⇒ The two-gap superconductivity disappears.

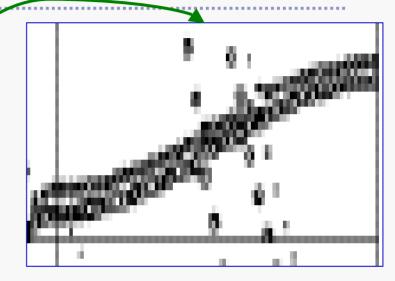
## 2x3 Supercell & CPA (x=1/12)



- C substitution
- 1. a change in the average potential in the B-C layer
- 2. the breaking of symmetry by C replacement of B in the supercell
- the splitting of the  $\sigma$  bands along  $\Gamma\text{-}A$  by 60 meV

## continued

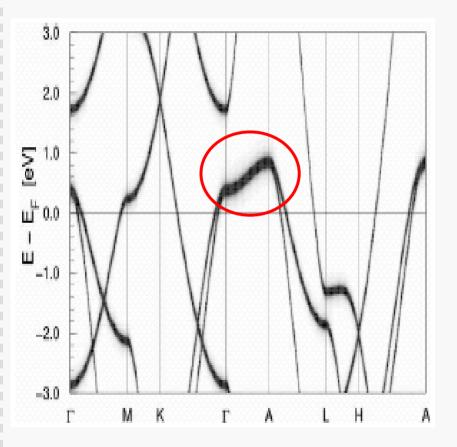




• Lower band is pulled down due to the stronger potential of the C atom than that of the B atom.

• The splitting provides an energy scale for  $\sigma$ -band broadening  $\gamma_0 \approx 7$  meV/%C.

# **CPAll (***x=0.0001***)**



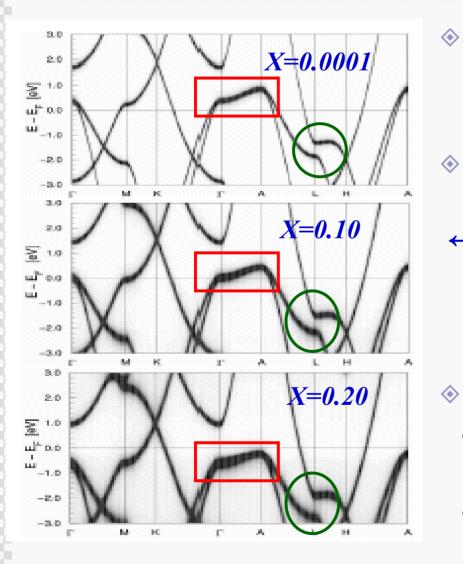
Since the numerical algorithms cannot reproduce the δ-function bands for  $x → \theta$ , it is included as a reference for the algorithmic contribution to the width.

### Width:

algorithmic contribution

- + disorder effect
- + (order effect)

# **CPA** III (Width of $\sigma$ band)



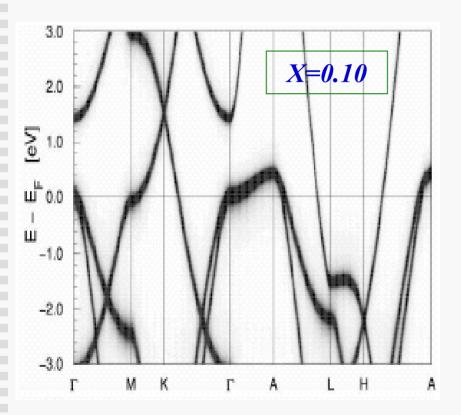
 $\delta E_F = -0.4 \text{ eV for } x = 0.10$ -1.0 eV for x = 0.20

- The band shift is not entirely rigid.
- The two valence bands at the L point split apart and broaden with increasing C.

 $\diamond$  averaged  $\gamma \approx 0.21 \text{ eV}$ 

- 0.06 eV from the B/C onsite energy difference
- 0.15 eV from disorder itself

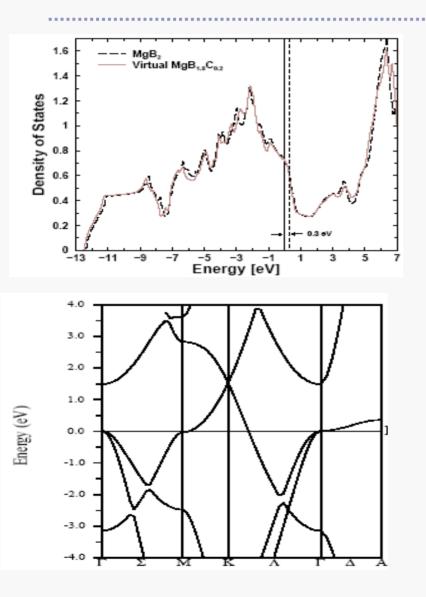
# **CPA IV (mean free path)**



The width corresponds to a  $\diamond$ width in wavevector given by  $\gamma = v_F \delta k$ .  $\therefore l_F = 2\pi v_F / \gamma$ At x=0.10,  $l_F$  may vary  $\diamond$ considerably over the FS due to the anisotropy of the  $\mathcal{V}_F$  .  $\diamond$  At  $\Gamma$ , the cylinder radius

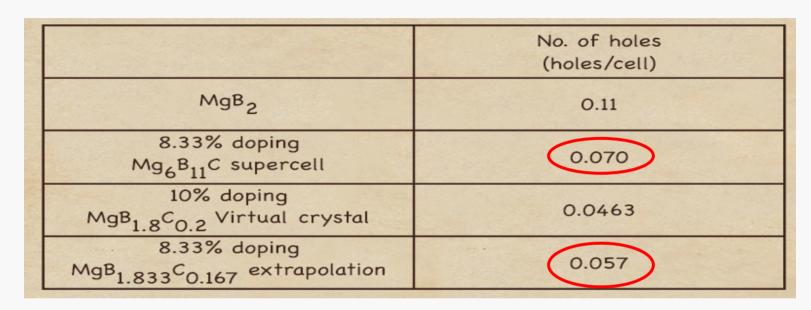
- shrunk to a point and the very small z-component of  $\mathcal{V}_F$ .
- ⇒ a very small  $l_F$  in the zdirection of the order of the layer spacing c.

# VCA (*x*=0.10)



- 1% increase in the occupied bandwidth
- The raising of E<sub>F</sub> by
   0.3 eV
- The cylindrical FS radii of the σ bands at Γ vanish.
- ⇒ corresponding to a topological transition

## $\sigma$ hole concentration



- Both the σ band FS are still intact, consistent with the 2-band superconductivity with substantial T<sub>c</sub> as seen in experiments.
- ♦ VCA is not reliable for substitutional C.
  - $\Rightarrow$  C cannot be thought simply as "B + e¯".

# Conclusion

- The  $\sigma$  band holes begin to disappear rapidly for x > 0.10.
- CPA reproduces the σ filling well, while VCA overestimates it.
- The largest disorder occurs in the 2s region in the lower valence band.
- As the σ band fills, there will be very strong deviation from "business as usual" in the coupled el.-ph. system.
   (The strength of coupling of bond-stretching modes with Q < 2k<sub>F</sub> continues to increase.)
- $Mg(B_{1-x}C_x)_2$  as well as  $Mg_{1-x}Al_xB_2$  is a good system to study the evolution of the unusual el.-ph. Coupling character.